

We also need estimates for the fundamental solutions  $\phi_{\pm}(x, k) = R_0(\lambda \pm i0)1 = \frac{1}{4} \left( \frac{k}{2\sqrt{|x|}} \right)^{\frac{n-2}{2}} H_{\frac{n-2}{2}}^{(1,2)}(k|x|)$   
 $(= \frac{e^{\pm ik|x|}}{4\sqrt{|x|}}, \text{ when } n=3)$ .

15.7 Lemma We have for all  $k > 0$  in  $\mathbb{R}^3$

$$(i) \quad \frac{e^{\pm ik|x-y|}}{|x-y|} = \frac{e^{\pm ikr}}{r} \left\{ e^{\mp ik\theta \cdot y} + O\left(\frac{1}{|x|}\right) \right\}$$

$$(ii) \quad \frac{\partial}{\partial y_j} \frac{e^{\pm ik|x-y|}}{|x-y|} = \frac{e^{\pm ikr}}{r} \left\{ \frac{\partial}{\partial y_j} e^{\mp ik\theta \cdot y} + O\left(\frac{1}{|x|}\right) \right\}$$

where  $r = |x|$  and  $\theta = \frac{x}{|x|}$ . The estimates hold uniformly w.r.t.  $y \in K \in \mathbb{R}^n$ .

Proof We have

$$|x-y|^2 = x^2 - 2x \cdot y + y^2 = |x|^2 \left( 1 - \frac{2\hat{x} \cdot y}{|x|} + O\left(\frac{1}{|x|^2}\right) \right)$$

$$\Rightarrow |x-y| = |x| \left( 1 - \frac{1}{2} \left( \frac{2\hat{x} \cdot y}{|x|} \right) + O\left(\frac{1}{|x|^2}\right) \right) \\ = |x| - \frac{\hat{x} \cdot y}{|x|} + O\left(\frac{1}{|x|^2}\right)$$

$\Rightarrow$  claim.  $\square$

Similar asymptotics hold for  $H_{\frac{n-2}{2}}^{(1,2)}(k|x|)$ .

Which yields

15.8 Lemma Assume  $f \in B$ . The asymptotics of  $u^\pm = R_0(\lambda \pm i0) f$  is given by

$$u^\pm(x) = \frac{C_n k^{\frac{n-3}{2}} e^{\pm ik|x|}}{|x|^{\frac{n-1}{2}}} \hat{f}(\pm k\theta) + u_0^\pm$$

where  $u_0^\pm \in B^*$ . Here  $\hat{f}(\xi) = \int_{\mathbb{R}^n} e^{-i\xi \cdot x} f(x) dx$ .  
 $C_n = \frac{1}{4\pi} \left( \frac{1}{2\pi i} \right)^{\frac{n-3}{2}}$

Proof  
 We prove the claim only in the case  $n=3$ .

$$u^\pm(x) = \int \frac{e^{\pm ik|x-y|}}{4\pi|x-y|} f(y) dy$$

$$\begin{aligned} & \stackrel{L. 17.7. (i)}{=} \frac{e^{\pm ik|x|}}{4\pi|x|} \int \left( \frac{e^{\mp ik\theta \cdot y}}{4\pi} + O\left(\frac{1}{|y|}\right) \right) f(y) dy \\ & = \frac{e^{\pm ik|x|}}{4\pi|x|} \hat{f}(\pm k\theta) + u_0^\pm \end{aligned}$$

Here

$$\|u_0^\pm(k\theta)\| \leq \frac{C}{(1+|k|)^2} \text{ for } f \in L^2_{comp}$$

Thus if

$$f_n = \chi_{B_n} f \in L^2_{comp}$$

we have

$$\|f_n - f\|_B \rightarrow 0 \text{ and}$$

$$f_n \in B^*$$

Thus

$$\begin{aligned} & \|u_n^\pm - u^\pm\|_{B^*} \leq \left\| \frac{e^{\pm ik|x|}}{|x|^{\frac{n-1}{2}}} \left( \hat{f}_n(\pm k\theta) - \hat{f}(\pm k\theta) \right) \right\|_{B^*} \\ & \rightarrow 0 \\ & \geq \|u_{n,0}^\pm - u_{0,\pm}^\pm\|_{B^*} \end{aligned}$$

By Theorem 14.1.1, [Hö] the map

$$f \mapsto \hat{f}(k\theta)$$

is bounded (isomorphism) from  $B \rightarrow L^2(\mathbb{T}^n)$ .

Thus

$$\int \left| \hat{f}_n(k\theta) - \hat{f}(k\theta) \right|^2 d\theta \rightarrow 0 \quad \text{as } n \rightarrow \infty$$

$$\Rightarrow \left\| \frac{e^{\pm i h x}}{h} \left( \hat{f}_n(k\theta) - \hat{f}(k\theta) \right) \right\|_{B^*} \rightarrow 0$$

$$\Rightarrow \left\| u_{n,0}^{\pm} - u_0^{\pm} \right\|_{B^*} \rightarrow 0$$

$$\Rightarrow u_0^{\pm} \in B^*.$$

□

Proof of Theorem 15.6:

$$(i) \quad u^0 \in B^0 \Leftrightarrow u^0 \in B^* \text{ and } (\Delta + h^2)u^0 = 0$$

$$\Rightarrow u = u^{\pm} = \int \tilde{e}^{\pm i h \theta \cdot x} v^{\pm}(\theta) d\theta + 0$$

Here

$$v^{\pm} \in L^2(S^{n-1})$$

For  $v_j^{\pm} \in C^{\infty}(S^{n-1})$  we have for

$$u_j = \int_{S^{n-1}} \tilde{e}^{\pm i h \theta \cdot x} v_j^{\pm}(\theta) d\theta$$

by (7.7.25) that

$$u_j(\pm\theta) = \left( \frac{2\pi h}{h} \right)^{\frac{n-1}{2}} \left[ v_j^{\pm}(k\theta) e^{-i h k} + v_j^{\pm}(-k\theta) e^{i h k} \right] + o(h^{-2})$$

Since  $\mathcal{O}(\frac{1}{r^2}) \in B^*$  and

$F: L^2(S^{n-1}) \rightarrow B^*$  is bounded

we get in the limit  $j \rightarrow \infty$

$$u^+(r\theta) = \left(\frac{2\pi k}{r}\right)^{\frac{n-1}{2}} \left[ v_+(r\theta) e^{ikr} + v_+(-r\theta) e^{-ikr} \right] + u^0$$

where  $u^0 \in B^*$ .

(ii) By definition,  $u^+ \in B^* \Leftrightarrow u \in B^*$  and

$$u^+ = R_0(\lambda + i0) f, \quad f \in B$$

Lemma 15.8  $\Rightarrow$

$$u^+ = C_n \frac{1}{k} \left(\frac{k}{r}\right)^{\frac{n-1}{2}} e^{ikr} \hat{f}(k\theta) + u_0 \in B^*$$

(iii) Since

$$u = u_0 + R(\lambda + i0) V u, \quad u_0 = \int e^{-i\theta \cdot x} v_+(\theta) d\theta$$

the claim follows from (i) and (ii).  $\square$

Remark

If

$$u = u_0 + R_0(\lambda + i0) V u$$

then

$$\mu^m(\theta) = (2\pi k)^{\frac{n-1}{2}} v_+(\theta)$$

when

$$u_0 = \int_{S^{n-1}} e^{-i\theta \cdot x} v_+(\theta) d\theta \quad \text{and}$$

$$\mu^m(\theta) = C_n k^{\frac{n-3}{2}} \int e^{-ix \cdot k\theta} V(x) u(x) dx$$

$S^m(\theta)$  is called the scattering amplitude.

We define three maps

$$h^0: B^0 \rightarrow L^2(S^{n-1})$$

$$h^m: B^m \rightarrow L^2(S^{n-1}) \quad \text{and}$$

$$h^+: B^+ \rightarrow L^2(S^{n-1})$$

by

$$b^0: u^0 \mapsto \nu^0$$

$$b^m: u^m \mapsto \nu^m \quad \text{and}$$

$$b^+: u^+ \mapsto \nu^+$$

15.9 Theorem  $b^0$  and  $b^m$  are

isomorphisms and  $b^+$  is surjective.

Moreover

$$\text{Ker } b^+ = \{ R(\lambda + i0) f \mid \hat{f}|_{\mathbb{R}S^{n-1}} = 0 \}$$

Proof Rellich  $\Rightarrow$   $b^0$  and  $b^m$  are injective

Since

$$u^0 = c \int_{\mathbb{R}S^{n-1}} e^{+i\lambda \cdot x} \nu_0(\theta) d\theta$$

the operator  $b^0$  is also surjective.

Since for any  $u^0$

$$u \in u_0 + R(\lambda + i0) \forall u$$

has a unique solution in  $B_{p_0}^*$ , we

conclude that  $b_1^m$  is also surjective.

How to prove  $b^+$  is surjective?

$$b^+ : U^+ \rightarrow V^+$$

$$u^+(x) = R_0(x+i0) f = \frac{e^{i\lambda x}}{\lambda^{n-1}} \mu_+(x) \in B^+$$

where

$$(1) \quad \mu_+(x) = c \hat{f}(x) \quad , \quad c = R^{\frac{n-2}{2}} C_1$$

$$f = \sqrt{u^m} \in B$$

Now claim follows from the fact that

$$f \mapsto \hat{f}(x) \text{ is surjective from } B \rightarrow L^2(S^{n-1})$$

(This is Theorem 14.1.1 [Hö])

Finally (2)  $\rightarrow$  (1)  $\square$

Recall

$$S : u^0 \mapsto u^m, \quad \begin{matrix} u^0 \in B^0 \\ u^m \in B^m \end{matrix}$$

$$\Sigma : v_- \mapsto v_+, \text{ where}$$

$$u^m = u_\pm + R(\lambda \mp i0) v u^m, \quad \hat{u}_\pm = v_\pm dS$$

Finally we have the operator

$$A : \mu^m \rightarrow f$$

The next lemma establishes the connection between them.

15.10 Lemma We have

(i)  $A = I - \Sigma$  and

(ii)  $\forall u^0, w^0 \in B$  with for fields  $\mu_0$  and  $\nu_0$  we have

$$\int_{\mathbb{R}^n} u^0 \mu \Sigma w^0 = -2ik \int_{S^{n-1}} \mu^0 A \nu^0, \quad C_n$$

Proof (i) By Lemma 15.8

$$R(\lambda \pm i0) f \sim C \frac{e^{\pm i|h|x|}}{|x|^{\frac{n-1}{2}}} \hat{f}(\pm h\theta)$$

Since

$$u = u_{\pm}^0 + R(\lambda \mp i0) v u$$

$$\hat{u}_{\pm}^0(k\theta) = v^{\pm}(k\theta) dS$$

T. 15.6 (i)

$\Rightarrow$

$$u_{\pm}^0 \sim v^{\pm}(k\theta) \frac{e^{ikh\tau}}{r^{\frac{n-1}{2}}} + v^{\pm}(\pm k\theta) \frac{e^{\mp ik\tau}}{r^{\frac{n-1}{2}}}$$

$\Rightarrow$

$$\begin{aligned} u^m &\sim (v^+(k\theta) + \hat{f}(k\theta)) \frac{e^{ikh\tau}}{r^{\frac{n-1}{2}}} + v^+(-k\theta) \frac{e^{-ikh\tau}}{r^{\frac{n-1}{2}}} \\ &\sim v^-(k\theta) \frac{e^{ikh\tau}}{r^{\frac{n-1}{2}}} + (v^+(-k\theta) + \hat{f}(k\theta)) \frac{e^{-ikh\tau}}{r^{\frac{n-1}{2}}} \end{aligned}$$

Thus  $\mu^m(\theta) = v^+(k\theta)$ ,  $\nu(\theta) = \hat{f}(k\theta)$

and

$$v^-(\theta) = \mu^m(\theta) + \gamma(\theta)$$

Thus

$$\Sigma \mu^m = \Sigma v^+ = v^- = \mu^m + \gamma = (I + A) \mu^m$$

which proves

$$\Sigma = I + A$$

(ii) By writing  $w^m = w^0 + w^+$  we have

$$\begin{aligned} \int_{\mathbb{R}^n} \bar{u}_0 b^2 m S w^0 &= \int_{\mathbb{R}^n} \bar{u}_0 b^2 m w^m \\ &= - \int_{\mathbb{R}^n} \bar{u}_0 (\Delta + h^2) w^m = - \int_{\mathbb{R}^n} \bar{u}_0 (\Delta + h^2) w^+ \\ &= \lim_{R \rightarrow \infty} - \int_{B_R} \bar{u}_0 (\Delta + h^2) w^+ = \\ &= \lim_{R \rightarrow \infty} \int_{\partial B_R} \frac{\partial \bar{u}_0}{\partial \nu} w^+ - \bar{u}_0 \frac{\partial w^+}{\partial \nu} = \lim_{R \rightarrow \infty} \int_{\partial B_R} \frac{-i k \bar{u}_0(\theta) A \nu^0(\theta)}{R^{n-1}} \\ &= -2ik \int N^0 A \nu^0 \end{aligned}$$



By Theorem 15.6



Define

$$S^+ : B^0 \rightarrow B^{0*}$$

by

$$S^+ : w^0 \mapsto -\frac{1}{2ik} \langle m S w^0, \cdot \rangle$$

### 15.11 Theorem

$$A = \mathcal{H}^* S^+ \mathcal{H}$$

where  $\mathcal{H}$  is the Helgoltz operator

$$u^0 = \mathcal{H} \mu_0 = \int_{S^{n-1}} e^{i u^0 \cdot x} \mu(\theta) d\theta$$

Proof

$$\langle \mathcal{H}^* S^+ \mathcal{H} v^0, \mu_0 \rangle = \langle S^+ \mathcal{H} v^0, \mathcal{H} \mu_0 \rangle$$

$$= -\frac{1}{2ik} \int_{\mathbb{R}^n} \overline{\mathcal{H} \mu_0} m S \mathcal{H} v^0$$

$$= -\frac{1}{2ik} \int_{\mathbb{R}^n} \overline{u_0} m S w^0 = \int_{S^{n-1}} \overline{\mu^0} A v^0$$

$$= \langle A v^0, \mu_0 \rangle$$

□

Remark

An element  $z \in B^{0*}$  can be interpreted

as follows:  $u^0 \in B^0 \Rightarrow u^0 = \mathcal{H} \mu_0$ . (Note)

$\mathcal{H} : L^2 \rightarrow B^0$  and  $\mathcal{H}^* : B^{0*} \rightarrow L^2$  and

by denoting  $v^0 = \mathcal{R}^* \ell$  we have for  
any  $\gamma^0 \in L^2(S^{n-1})$

$$\langle v^0, \gamma^0 \rangle_{S^{n-1}} = \langle \mathcal{R}^* \ell, \gamma^0 \rangle = \langle \ell, \mathcal{R} \gamma^0 \rangle$$

Thus

$$\langle \ell, u^0 \rangle = \langle v^0, \gamma^0 \rangle, \quad \forall u^0 = \mathcal{R} \gamma^0.$$

16. A Generalized Scattering Operator and interior transmission problems  
In Section 15 real and sources

$$u^0 = \mathcal{R} p_0$$

to illuminate the scatterer  $m$ .

Other sources might be supported in a compact set  $D$ .

Let  $m \in L^\infty$ ,  $\text{supp } m \subset D$ , and  $u^0 \in L^2(D)$ .

By Theorem 15.3  $\exists! u^+ \in B^+$  solving

$$(II) \quad (\Delta + k^2(1+m)) u^+ = k^2 m u^0$$

$$\Rightarrow u^m = u^0 + u^+ \in L^2(D)$$

Note that  $u^+$  is outgoing in the whole  $\mathbb{R}^n$ .

v) want to extend the scattering map  $S$  to such sources.

Let

$$B_D^0 = \{ w \in L^2(D) \mid (\Delta + V)w = 0 \}$$

$$B_D^m = \{ w \in L^2(D) \mid (\Delta + V + m)w = 0 \}$$

We define

$$S_D : B_D^0 \rightarrow B_D^m, \quad u^0 \mapsto u^m \quad \text{and}$$

$$S_D^+ : B_D^0 \rightarrow B_D^{0*}$$

$$S_D^+ : w^0 \mapsto \frac{1}{2i} \langle m, S_D w^0, \cdot \rangle$$

Clearly, since the right hand side is independent how  $u^0$  is defined outside of  $\text{supp } m$  we have for  $u^0, w^0 \in B^0$

$$S_D u^0 = S u^0|_D$$

$$\langle S_D^+ u^0, w^0 \rangle = \langle S^+ u^0, w^0 \rangle$$

The interior transmission problem (ITP)

From now on denote

$$P^0 = (\Delta + k^2)$$

$$P^m = (\Delta + k(1+m))$$

$$u \in H^k(D) \Leftrightarrow \|\partial^\alpha u\|_{L^2} < \infty \quad \forall |\alpha| \leq k$$

$\dot{H}^k(D)$  is the closure of  $C_c^\infty$  in  $H^k(D)$

16.1. Definition  $k$  is a  $D$ -transmission

eigenvalue for  $m \in L^\infty(D)$ , if  $\exists$

$$u^0 \in B_D^0 \text{ and } u^m \in B_D^m \text{ s.t.}$$

$$0 \neq u^0 - u^m \in \dot{H}^2(D)$$

We also say that  $k$  is a transmission

eigenvalue (TE) if  $\exists u_0 \neq 0$  s.t. the

scattered field is both incoming and

outgoing

$$u^m = u^0 + \underbrace{R(\lambda + i0)}_{=: \mathcal{R}^m} \mathcal{V} u^m$$

$$\mathcal{V} = \frac{1}{k} \Delta u^m$$

Since  $\text{supp } \mathcal{V}$  is compact we have

$u^s \equiv 0$  outside  $\text{supp } \mathcal{V}$

Now if  $\text{supp } v \subset D$  we have

$$u^0 \in B_D^0, \quad u^m \in B_D^m$$

and

$$u^0 - u^m = u^s$$

vanishes in the neighborhood of  $\partial D$ .

Thus  $u^0 - u^m \in H_0^2(D)$  and the

interior transmission problem

$$\begin{aligned} (\Delta + L^2)v &= 0 & \text{in } D \\ (\Delta + L^2(l+m))w &= 0 & \text{in } D \end{aligned}$$

$$v = w, \quad \frac{\partial v}{\partial \nu} = \frac{\partial w}{\partial \nu} \quad \text{on } \partial D$$

has a non-trivial solution  $v = u^0$

and  $w = u^m$ .